

# Second-order hyperbolic Fuchsian systems

## Applications to Einstein vacuum spacetimes

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Fuchsian systems:

$$Du(t, x) + N(x)u(t, x) = tf(t, x, u, \partial_x u, Du),$$

$$\lim_{t \searrow 0} u(t, x) = 0.$$

with  $D := t\partial_t$ .

## What does the Fuchsian theory provide?

Theory of **singular PDEs** which allows to construct solutions with a **prescribed singular behavior**.

Applications in General Relativity:

- Gowdy vacuum solutions: Kichenassamy-Rendall, Rendall.
- Polarized  $T^2$ -symmetric vacuum solutions: Isenberg-Kichenassamy.
- Cosmological models with stiff fluid/scalar field: Andersson-Rendall.
- ...

# Some drawbacks of the current theory

- Often not Fuchsian equations directly. First need to identify the “leading order part” of the unknowns at the singularity.
- Often second-order equations.
- Numerical scheme for **hyperbolic** Fuchsian equations and proof of convergence?

## Aims of the talk today:

- Tackle these drawbacks for a subset of Fuchsian equations.
- Present numerical examples.

## Class of equations

$$(D \circ D)u(t, x) + 2a(x) Du(t, x) + b(x) u(t, x) \\ = f(t, x, u, \partial_x u, Du, \partial_x^2 u, \partial_x Du),$$

with

- 1  $u : ]0, \delta] \times \mathbb{R} \rightarrow \mathbb{R}$ , periodic in space for a  $\delta > 0$ ,
- 2  $a, b$  are smooth periodic functions on  $\mathbb{R}$ ,
- 3  $f$  is the source with certain properties, see below.

Direct generalizations:

- Systems.
- More than one spatial dimensions.

## Canonical two-term expansion

$$u(t, x) = \begin{cases} u_*(x)t^{-\lambda_1} + u_{**}(x)t^{-\lambda_2} + O(t^{-\lambda_2+\alpha}), & a^2(x) > b(x), \\ u_*(x)t^{-\lambda_1} \log t + u_{**}(x)t^{-\lambda_1} + O(t^{-\lambda_1+\alpha}), & a^2(x) = b(x), \end{cases}$$

at  $t = 0$  for  $\alpha > 0$ , where

$$\lambda_1(x) := a(x) + \sqrt{a^2(x) - b(x)}, \quad \lambda_2(x) := a(x) - \sqrt{a^2(x) - b(x)}.$$

## Singular initial value problem (SIVP)

Does there exist a unique solution of the Fuchsian system in a given regularity class which obeys the canonical two-term expansion at  $t = 0$  for *given asymptotic data*  $u_*(x)$ ,  $u_{**}(x)$ ?

## Class of hyperbolic equations

Consider

$$\begin{aligned}(D \circ D)u(t, x) + 2a(x) Du(t, x) + b(x) u(t, x) \\ = f(t, x, u, \partial_x u, Du) + t^2 c^2(t, x) \partial_x^2 u(t, x)\end{aligned}$$

with the speed of propagation

$$c(t, x) = t^{\beta(x)} k(t, x).$$

Here,

- $\beta(x)$  is a smooth periodic function larger than  $-1$ .
- $k(t, x)$  is a smooth positive spatially periodic function so that all derivatives have finite limits at  $t = 0$ .

# Examples of 2nd-order hyperbolic Fuchsian equations

## 1 Euler–Poisson–Darboux equation:

$$(D \circ D)u - \Delta\lambda Du - t^2 \partial_x^2 u = 0,$$

for  $\Delta\lambda \geq 0$ . Equivalently  $\partial_t^2 u - \partial_x^2 u = \frac{1}{t}(\Delta\lambda - 1)\partial_t u$ .

SIVP: Look for solutions of the form

$$u(t, x) = \begin{cases} u_*(x) + u_{**}(x)t^{\Delta\lambda} + \dots & \Delta\lambda > 0, \\ u_*(x) \log t + u_{**}(x) + \dots & \Delta\lambda = 0, \end{cases}$$

## 2 (Main evolution part of the) Gowdy vacuum equations:

$$D^2 P - t^2 \partial_x^2 P = e^{2P} (DQ)^2 - t^2 e^{2P} (\partial_x Q)^2,$$

$$D^2 Q - 2k DQ - t^2 \partial_x^2 Q = -2(k + DP)DQ + 2t^2 \partial_x P \partial_x Q.$$

SIVP: Look for solutions of the form

$$P(t, x) = -k(x) \log t + P_{**}(x) + \dots,$$

$$Q(t, x) = Q_*(x) + Q_{**}(x)t^{2k(x)} + \dots$$

## The space $X_{\delta,\alpha,1}$

For a given second-order Fuchsian equation and some  $\delta, \alpha > 0$ , let  $X_{\delta,\alpha,1}$  be the space of functions

$$w \in C^0((0, \delta], H^1(U)) \cap C^1((0, \delta], L^2(U)) \text{ and } \|w\|_{\delta,\alpha,1} < \infty,$$

with the standard norm  $\|\cdot\|_{\delta,\alpha,1}$  **weighted by the factor  $t^{\lambda_2 - \alpha}$** .

## SIVP

We say that  $u$  is a solution of the SIVP for given asymptotic data  $u_*, u_{**} \in H^1(U)$ , if  $u$  is a weak solution of the 2nd-order hyperbolic Fuchsian eq. and there exists  $w \in X_{\delta,\alpha,1}$  so that

$$u(t, x) = u_{\text{canonic}}(t, x) + w(t, x).$$

The function  $w$  is called **remainder**.



# Well-Posedness in an important case

Consider the case

$$f(t, x, u, \partial_x u, Du) = f_0(t, x)$$

for a given function  $f_0$ .

## Result

For any asymptotic data  $u_*, u_{**} \in H^2(U)$ , there exists a unique solution of the SIVP with remainder  $w \in X_{\delta, \alpha, 1}$  provided:

- 1 we can choose  $\delta, \alpha > 0$  so that the matrix

$$\begin{pmatrix} \lambda_1 - \lambda_2 + \alpha & -\eta/2 & 0 \\ -\eta/2 & \alpha & t\partial_x c - \partial_x(\lambda_1 - \lambda_2)(tc \ln t) \\ 0 & t\partial_x c - \partial_x(\lambda_1 - \lambda_2)(tc \ln t) & \lambda_1 - \lambda_2 + \alpha - 1 - Dc/c \end{pmatrix}$$

is positive semidef. at each  $(t, x) \in (0, \delta) \times U$  for a  $\eta > 0$ .

- 2  $f_0 \in X_{\delta, \alpha + \varepsilon, 0}$  for some  $\varepsilon > 0$ .
- 3  $\alpha < 2(\beta(x) + 1) - (\lambda_1(x) - \lambda_2(x))$  for all  $x \in U$ .

# Main ingredients of the existence proof

- 1 **Main idea:** Solve a sequence of **regular initial value problems**, i.e. a sequence of approximate solutions  $(u_n)$  with initial times  $(\tau_n)$  going to zero as  $n \rightarrow \infty$ .
- 2 Obtain “good” energy estimates consistent with the space  $X_{\delta,\alpha,1}$  if the energy dissipation matrix is positive semidefinite.
- 3 Convergence:

$$\|u_n - u_m\|_{\delta,\alpha,1} \leq C|G(\tau_n) - G(\tau_m)|,$$

where

$$G(t) := \int_0^t s^{-1} \|s^{\lambda_2 - \alpha} (f_0(s) - L[u_{\text{canonic}}](s))\|_{L^2(U)} ds,$$

and  $C > 0$  independent of  $n$ . Under the hypothesis, we have  $\lim_{t \rightarrow 0} G(t) = 0$ . Hence  $(w_n)$  is a Cauchy sequence in  $X_{\delta,\alpha,1}$ .

- 4 Check that the limit function  $u$  satisfies the weak equation.

## Regular initial value problem (RIVP)

We say that  $\tilde{u}$  is a solution of the RIVP with **initial time**  $\tau_0 > 0$ , provided (after mollification)

For  $t \in (0, \tau_0)$  :

$$\tilde{u}(t, \mathbf{x}) = u_{\text{canonic}}(t, \mathbf{x}).$$

For  $t \in [\tau_0, \delta)$  :  $\tilde{u}$  is a smooth solution of the 2nd-order Fuchsian hyperbolic equation and initial data

$$\begin{aligned}\tilde{u}(\tau_0, \mathbf{x}) &= u_{\text{canonic}}(\tau_0, \mathbf{x}), \\ \partial_t \tilde{u}(\tau_0, \mathbf{x}) &= \partial_t u_{\text{canonic}}(\tau_0, \mathbf{x}).\end{aligned}$$

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## Comments on this result:

- Main difference to standard Fuchsian theory: approximate by smooth solutions of RIVP with an explicit convergence statement in the norm. **Leads to a numerical scheme!**
- Loss of regularity.

## Generalizations of the rigorous results:

- 1 General (non-linear) sources:

**Idea:** Iterate over the special case before.

**Result:** If source term is locally Lipschitz in the space  $X_{\delta,\alpha,1}$  and  $\delta > 0$  is sufficiently small, then obtain a fixed point iteration. Strong convergence in the norm!

- 2 Increase regularity assumptions to control arbitrarily many derivatives in spaces  $X_{\delta,\alpha,k}$  with  $k \in \mathbb{N}$ .

Consider the case

$$f(t, x, u, \partial_x u, Du) = f_0(t, x)$$

for a given function  $f_0$ .

## Result

For any asymptotic data  $u_*, u_{**} \in H^2(U)$ , there exists a unique solution of the SIVP with remainder  $w \in X_{\delta, \alpha, 1}$  provided:

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# Example application: Euler-Poisson-Darboux eq.

- Consider

$$(D \circ D)u - \Delta \lambda Du - t^2 \partial_x^2 u = 0,$$

for a constant  $\Delta \lambda \geq 0$ , and find solutions of the form

$$u(t, x) = \begin{cases} u_*(x) + u_{**}(x)t^{\Delta \lambda} + \dots & \Delta \lambda > 0, \\ u_*(x) \log t + u_{**}(x) + \dots & \Delta \lambda = 0. \end{cases}$$

- It turns out: well-posedness theory applies for  $0 \leq \Delta \lambda < 2$ .
- Periodicity in space  $\rightarrow$  explicit solutions in terms of Bessel functions. Result:

**For  $0 \leq \Delta \lambda < 2$ :** Solutions are consistent with the canonical two-term expansion.

**For  $\Delta \lambda = 2$ :** Solutions behave like

$$u(t, x) = u_*(x) + u_{**}(x)t^{\Delta \lambda} \log t + \dots$$

Second spatial derivative becomes significant at  $t = 0$ !



General numerical approach, using EPD-equation as a model:

- 1 Introduce time variable  $\tau := \ln t$ . Then equation becomes

$$\partial_\tau^2 u - \Delta \lambda \partial_\tau u - e^{2\tau} \partial_x^2 u = 0.$$

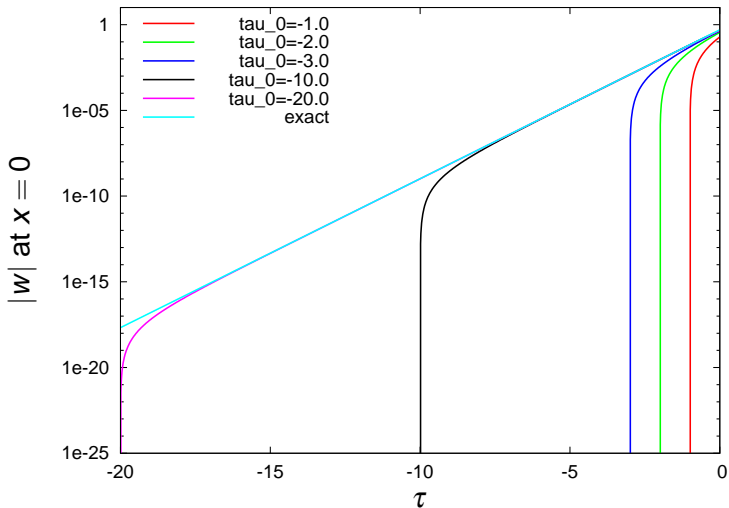
Singularity is “shifted to  $\tau = -\infty$ ”.

- 2 Write the equation for the remainder  $w$ , after having fixed asymptotic data  $u_*$ ,  $u_{**}$ .
- 3 Solve sequence  $(w_n)$  of solutions of RIVPs with initial times  $\tau_n \rightarrow -\infty$  with data

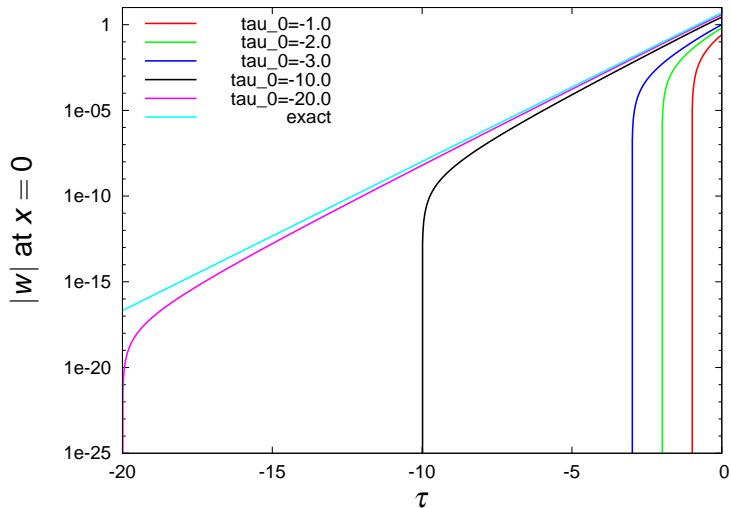
$$w_n(\tau_n) = 0, \quad \partial_\tau w_n(\tau_n) = 0,$$

using a direct discretization of the second-order equation (as suggested by Kreiss).

Choice:  $\Delta\lambda = 1.0$ ,  $u_* = \cos x$ ,  $u_{**} = 0$ .



Choice:  $\Delta\lambda = 1.9$ ,  $u_* = \cos x$ ,  $u_{**} = 0$ .



## Desired generalizations of the current theory:

- Mixed first- and second-order hyperbolic Fuchsian systems.

## Some ideas for applications:

- Gowdy coupled to a perfect fluid. Interaction of cosmological singularity and shocks?
- Numerical construction and analysis of solutions with Cauchy horizons or “pieces of Cauchy horizons” in the Gowdy case?