

A posteriori analysis of a discontinuous Galerkin scheme for a diffuse interface model

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Outline

1. Van-der-Waals fluid model
2. Relative entropy stability framework
 - convex energy
 - non-convex energy
3. Semi-discrete dG scheme
4. Reconstruction and error estimate

Van-der-Waals fluid model

One space dimension; describing fluid flows undergoing liquid-vapor phase transition:

$$\begin{aligned}u_t - v_x &= 0 \\v_t - W'(u)_x &= \mu v_{xx} - \gamma u_{xxx}\end{aligned}\tag{vdW}$$

- u specific volume, v velocity
- W **non-convex** energy density \Rightarrow (vdW) is hyperbolic-elliptic
- $W \in C^3(\mathbb{R}, [0, \infty))$
- $\gamma > 0$ capillarity parameter, $\mu \geq 0$ viscosity.

We consider the problem on the flat circle, denoted S^1 .

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Associated energy balance

$$\left(W(u) + \frac{\gamma}{2}(u_x)^2 + \frac{1}{2}v^2 \right)_t - \left(vW'(u) - \gamma v u_{xx} + \gamma v_x u_x + \mu v v_x \right)_x + \mu (v_x)^2 = 0.$$

Recall: Standard relative entropy (W strictly convex)

Consider the non-regularized, hyperbolic problem

$$\begin{aligned}u_t - v_x &= 0 \\v_t - W'(u)_x &= 0\end{aligned}$$

with W **strictly convex**.

Solutions (u, v) , (\tilde{u}, \tilde{v}) can be compared by their relative entropy

$$\int_{S^1} W(\tilde{u}) - W(u) - W'(u)(\tilde{u} - u) + \frac{1}{2}(\tilde{v} - v)^2 \, dx,$$

which is equivalent to

$$\|\tilde{u} - u\|_{L^2(S^1)}^2 + \|\tilde{v} - v\|_{L^2(S^1)}^2.$$

Recall: Standard relative entropy (W strictly convex)

The relative entropy satisfies

$$\begin{aligned} & \frac{d}{dt} \int_{S^1} W(\tilde{u}) - W(u) - W'(u)(\tilde{u} - u) + \frac{1}{2}(\tilde{v} - v)^2 dx \\ & - \int_{S^1} \left(\tilde{v} W'(\tilde{u}) - v W'(u) - W'(u)(\tilde{v} - v) - v(W'(\tilde{u}) - W'(u)) \right) dx \\ & \leq \int_{S^1} v_x (W'(\tilde{u}) - W'(u) - W''(u)(\tilde{u} - u)) dx. \end{aligned}$$

Standard relative entropy argument (W strictly convex)

For $|W'''|$ bounded and v Lipschitz this implies

$$\begin{aligned} & \frac{d}{dt} \int_{S^1} W(\tilde{u}) - W(u) - W'(u)(\tilde{u} - u) + \frac{1}{2}(\tilde{v} - v)^2 dx \\ & \leq C(\|\tilde{u} - u\|_{L^2(S^1)}^2 + \|\tilde{v} - v\|_{L^2(S^1)}^2) \\ & \leq C \int_{S^1} W(\tilde{u}) - W(u) - W'(u)(\tilde{u} - u) + \frac{1}{2}(\tilde{v} - v)^2 dx. \end{aligned}$$

Therefore, by Gronwall's Lemma, the relative entropy grows at most exponentially and, therefore, for $t > s$

$$\begin{aligned} & \|\tilde{u}(t, \cdot) - u(t, \cdot)\|_{L^2} + \|\tilde{v}(t, \cdot) - v(t, \cdot)\|_{L^2} \\ & \leq e^{C(t-s)} \left(\|\tilde{u}(s, \cdot) - u(s, \cdot)\|_{L^2} + \|\tilde{v}(s, \cdot) - v(s, \cdot)\|_{L^2} \right). \end{aligned}$$

Back to the multi-phase case, i.e., W not convex

Regularized model:

$$\begin{aligned}u_t - v_x &= 0 \\v_t - W'(u)_x &= \mu v_{xx} - \gamma u_{xxx}.\end{aligned}$$

For different (u, v) (\tilde{u}, \tilde{v}) their relative entropy is given by

$$\int_{S^1} W(\tilde{u}) - W(u) - W'(u)(\tilde{u} - u) + \frac{1}{2}(\tilde{v} - v)^2 + \frac{\gamma}{2}(\tilde{u}_x - u_x)^2 \, dx,$$

which is not convex, as W is not convex and γ is small.

On S^1 there are no boundary terms, and we obtain

$$\begin{aligned}& \frac{d}{dt} \int_{S^1} W(\tilde{u}) - W(u) - W'(u)(\tilde{u} - u) + \frac{1}{2}(\tilde{v} - v)^2 + \frac{\gamma}{2}(\tilde{u}_x - u_x)^2 \, dx \\& \leq \int_{S^1} v_x(W'(\tilde{u}) - W'(u) - W''(u)(\tilde{u} - u)) - \mu(v_x - \tilde{v}_x)^2 \, dx.\end{aligned}$$

Estimating the partial relative entropy rate

Idea: Remove the W terms from the relative entropy. Indeed

$$\begin{aligned} & \partial_t \left(W(\tilde{u}) - W(u) - W'(u)(\tilde{u} - u) \right) \\ &= W'(\tilde{u})\tilde{u}_t - W'(u)u_t - W''(u)u_t(\tilde{u} - u) - W'(u)\tilde{u}_t + W'(u)u_t \\ &= \tilde{v}_x W'(\tilde{u}) - \tilde{v}_x W'(u) - v_x W''(u)(\tilde{u} - u). \end{aligned}$$

Thus, we can shift the W terms to the right hand side and find

$$\begin{aligned} \frac{d}{dt} \int_{S^1} \frac{1}{2} (\tilde{v} - v)^2 + \frac{\gamma}{2} (\tilde{u}_x - u_x)^2 dx \\ \leq \int_{S^1} (v - \tilde{v})_x (W'(\tilde{u}) - W'(u)) - \mu (v_x - \tilde{v}_x)^2 dx. \end{aligned}$$

Continuous dependence on initial data

Lemma (JG '13)

For $\mu > 0$ let $u_0, \tilde{u}_0 \in H^3(S^1)$, $v_0, \tilde{v}_0 \in H^2(S^1)$ be given with $\int_{S^1} (u_0 - \tilde{u}_0) dx = 0$.

Then, for any $T > 0$, it exists a constant $C = C(u_0, v_0, \gamma, \mu, T)$ such that

$$\begin{aligned} \|v(t, \cdot) - \tilde{v}(t, \cdot)\|_{L^2(S^1)} + |u(t, \cdot) - \tilde{u}(t, \cdot)|_{H^1(S^1)} \\ \leq C \left(\|v_0 - \tilde{v}_0\|_{L^2(S^1)} + |u_0 - \tilde{u}_0|_{H^1(S^1)} \right) \end{aligned}$$

for all $t \leq T$.

Sketch of the proof: For any $T > 0$ there exist strong solutions

$$\begin{aligned} u, \tilde{u} &\in C^1((0, T), L^2(S^1)) \cap C^0((0, T), H^3(S^1)) \\ v, \tilde{v} &\in C^1((0, T), L^2(S^1)) \cap C^0((0, T), H^2(S^1)). \end{aligned}$$

Sketch of the proof

Due to the energy inequality and the continuous embedding

$H^1(S^1) \rightarrow C^0(S^1)$, we find that $\|u\|_{L^\infty}, \|\tilde{u}\|_{L^\infty} < \infty$.

The partial relative entropy calculation leads to

$$\begin{aligned} \frac{d}{dt} \int_{S^1} \frac{1}{2} (\tilde{v} - v)^2 + \frac{\gamma}{2} (\tilde{u}_x - u_x)^2 dx \\ = \int_{S^1} (\tilde{v} - v)_x (W'(u) - W'(\tilde{u})) - \mu (v_x - \tilde{v}_x)^2 dx. \end{aligned}$$

Such that by Young's and Poincaré's inequality

$$\frac{d}{dt} \int_{S^1} \frac{1}{2} (\tilde{v} - v)^2 + \frac{\gamma}{2} (\tilde{u}_x - u_x)^2 dx \leq C \int_{S^1} \frac{\gamma}{2} (\tilde{u}_x - u_x)^2 dx.$$

Gronwall's Lemma concludes the proof.

Continuous dependence on initial data for $\mu = 0$

Lemma (JG '13)

For $\mu = 0$ let solutions

$$u, \tilde{u} \in C^1((0, T), L^2(S^1)) \cap C^0((0, T), H^3(S^1))$$

$$v, \tilde{v} \in C^1((0, T), L^2(S^1)) \cap C^0((0, T), H^1(S^1))$$

corresponding to initial data (u_0, v_0) and $(\tilde{u}_0, \tilde{v}_0)$ be given with $\int_{S^1} (u_0 - \tilde{u}_0) dx = 0$. Let $|W''|$ be bounded.

Then it exists a constant $C = C(u_0, v_0, \gamma, T)$ such that

$$\begin{aligned} \|v(t, \cdot) - \tilde{v}(t, \cdot)\|_{L^2(S^1)} + |u(t, \cdot) - \tilde{u}(t, \cdot)|_{H^1(S^1)} \\ \leq C \left(\|v_0 - \tilde{v}_0\|_{L^2(S^1)} + |u_0 - \tilde{u}_0|_{H^1(S^1)} \right) \end{aligned}$$

for all $t \leq T$.

Sketch of the proof

The partial relative entropy calculation leads to

$$\begin{aligned} \frac{d}{dt} \int_{S^1} \frac{1}{2} (\tilde{v} - v)^2 + \frac{\gamma}{2} (\tilde{u}_x - u_x)^2 dx \\ &= \int_{S^1} (\tilde{v} - v)_x (W'(u) - W'(\tilde{u})) dx \\ &\leq \int_{S^1} (\tilde{v} - v)^2 + ((W'(u) - W'(\tilde{u}))_x)^2 dx \\ &\leq C \int_{S^1} \frac{1}{2} (\tilde{v} - v)^2 + \frac{\gamma}{2} (\tilde{u}_x - u_x)^2 dx \end{aligned}$$

using Young's and Poincaré's inequality.

Applying Gronwall's Lemma concludes the proof.

Semi-discrete dG scheme (from now on $\mu = 0$)

Decompose $[0, 1]$ into $0 = x_0 < x_1 < \dots < x_N = 1$. Identify 0 and 1.
 $\mathbb{V}_q :=$ space of (discontinuous) piece-wise polynomials of degree $\leq q$,
 $\mathbb{V}_q^c := \mathbb{V}_q \cap C^0(S^1)$.

Find $u_h, v_h \in C^1((0, T), \mathbb{V}_q)$, $\tau_h \in C^0([0, T], \mathbb{V}_q)$ such that

$$0 = \int_{S^1} \partial_t u_h \Phi - G[u_h] \Phi \, dx \quad \forall \Phi \in \mathbb{V}_q$$

$$0 = \int_{S^1} \partial_t v_h \Psi - G[\tau_h] \Psi \, dx \quad \forall \Psi \in \mathbb{V}_q$$

$$0 = \int_{S^1} \tau_h Z - W'(u_h) Z \, dx + \gamma a_h(u_h, Z) \quad \forall Z \in \mathbb{V}_q,$$

where a_h is an interior penalty discretisation of the Laplacian and G denotes a discrete gradient operator.

Discretisations

For $g_h \in \mathbb{V}_q$ we define $G[g_h] \in \mathbb{V}_q$ by

$$\int_{S^1} G[g_h] \Phi \, dx = \sum_{i=0}^{N-1} \left(\int_{x_i}^{x_{i+1}} (g_h)_x \Phi \, dx - \llbracket g_h \rrbracket_{x_i} \{ \Phi \}_{x_i} \right)$$

for all $\Phi \in \mathbb{V}_q$.

For $f_h, g_h \in \mathbb{V}_q$ we define

$$a_h(f_h, g_h) := \sum_{i=0}^{N-1} \left(\int_{x_i}^{x_{i+1}} (f_h)_x (g_h)_x \, dx - \llbracket f_h \rrbracket_{x_i} \{ (g_h)_x \}_{x_i} \right. \\ \left. - \llbracket g_h \rrbracket_{x_i} \{ (f_h)_x \}_{x_i} + \frac{\sigma}{h} \llbracket f_h \rrbracket_{x_i} \llbracket g_h \rrbracket_{x_i} \right),$$

for some $\sigma \gg 1$, such that $a_h : \mathbb{V}_q \times \mathbb{V}_q \rightarrow \mathbb{R}$ is coercive.

Reconstruction 1

$\hat{\tau} \in \mathbb{V}_{q+1}$ is defined by

$$0 = \int_{S^1} (\hat{\tau}_x - G[\tau_h]) \Psi \, dx \quad \forall \Psi \in \mathbb{V}_q \quad \text{and}$$
$$\hat{\tau}(x_n^-) = \frac{\tau_h(x_n^-) + \tau_h(x_n^+)}{2} \quad \forall n \in \{0, \dots, N-1\}.$$

It can be shown that

- $\hat{\tau}$ is continuous.
- $\|\hat{\tau} - \tau_h\|_{L^2(S^1)}^2 \leq \sum_{i=0}^{N-1} (x_{i+1} - x_i) \left(\llbracket \tau_h \rrbracket_{x_i}^2 + \llbracket \tau_h \rrbracket_{x_{i+1}}^2 \right).$

Reconstruction 2

$\tilde{u} \in H^3(S^1)$ is defined by

$$0 = \gamma \partial_{xx} \tilde{u} - P_{q+1}^c(W'(u_h)) + \hat{\tau},$$

where $P_{q+1}^c : L^2(S^1) \rightarrow \mathbb{V}_{q+1}^c$ is the L^2 -projection, and

$$\int_{S^1} (u - \tilde{u}) dx = 0.$$

Using elliptic regularity and a posteriori control of elliptic reconstructions we have

$$\begin{aligned} \|\tilde{u} - u_h\|_{dG} &:= \left(\sum | \tilde{u} - u_h |_{H^1((x_i, x_{i+1}))}^2 + \sum | [\tilde{u} - u_h]_{x_i} |^2 \right)^{\frac{1}{2}} \\ &\leq C \| P_{q+1}^c(W'(u_h)) - W'(u_h) \|_{H^{-1}(S^1)} \\ &\quad + C \| \hat{\tau} - \tau_h \|_{H^{-1}(S^1)} + H[u_h, W'(u_h), \tau_h], \end{aligned}$$

where $H[u_h, W'(u_h), \tau_h]$ is an explicitly computable estimator, expected to be of order h^q .

Reconstruction 2 cont'd

$$\begin{aligned} \|P_{q+1}^c(W'(u_h)) - W'(u_h)\|_{H^{-1}(S^1)} &\leq C \sqrt{\sum_i (x_{i+1} - x_i) |[[u_h]]|^2} \\ &\quad + C \sup_i (x_{i+1} - x_i)^{q+1} |W'(u_h)|_{H^1((x_i, x_{i+1}))}. \end{aligned}$$

According to arguments provided by Nochetto and Makridakis '06 terms of the structure

$$\sqrt{\sum_i (x_{i+1} - x_i) |[[\cdot]]|^2}$$

are expected to be of optimal order.

Reconstruction 3

$\hat{v} \in \mathbb{V}_{q+1}$ is defined by

$$0 = \int_{S^1} (\hat{v}_x - G[v_h]) \Psi \, dx \quad \forall \Psi \in \mathbb{V}_q \quad \text{and} \quad \hat{v}(x_n^-) = \frac{v_h(x_n^-) + v_h(x_n^+)}{2} \quad \forall n.$$

$\tilde{v} \in H^2(S^1)$ is defined by

$$0 = \partial_{xx} \tilde{v} - \partial_{xt} \tilde{u} \quad \text{and} \quad \int_{S^1} \tilde{v} - v_h \, dx = 0.$$

\hat{v} is continuous and

$$\|\hat{v} - v_h\|_{L^2(S^1)}^2 \leq \sum_{i=0}^{N-1} (x_{i+1} - x_i) \left(\llbracket v_h \rrbracket_{x_i}^2 + \llbracket v_h \rrbracket_{x_{i+1}}^2 \right),$$

while

$$\|\tilde{v} - \hat{v}\|_{L^2(S^1)} \leq \|\partial_t \tilde{u} - \partial_t u_h\|_{dG}.$$

Perturbed equation

By definition the reconstructions satisfy (point-wise a.e.)

$$\begin{aligned}\tilde{u}_t - \tilde{v}_x &= 0 \\ (v_h)_t - \hat{\tau}_x &= 0\end{aligned}$$

which implies, by definition of \tilde{u} ,

$$(v_h)_t - (P_{q+1}^c(W'(u_h)))_x + \gamma \tilde{u}_{xxx} = 0.$$

Thus,

$$\begin{aligned}\tilde{u}_t - \tilde{v}_x &= 0 \\ \tilde{v}_t - W'(\tilde{u})_x &= -\gamma \tilde{u}_{xxx} + E,\end{aligned}$$

where the residual E is given by

$$E := \partial_t(\tilde{v} - v_h) + \partial_x(P_{q+1}^c(W'(u_h)) - W'(\tilde{u})).$$

By the estimates above, E is bounded explicitly in terms of u_h, v_h, τ_h .

Partial relative entropy

Then, the partial relative entropy calculation implies

$$\begin{aligned} \frac{d}{dt} \int_{S^1} \frac{1}{2} (\tilde{v} - v)^2 + \frac{\gamma}{2} (\tilde{u}_x - u_x)^2 dx \\ = \int_{S^1} (v - \tilde{v})(W'(\tilde{u}) - W'(u))_x dx + \int_{S^1} (v - \tilde{v})E dx. \end{aligned}$$

We infer

$$\begin{aligned} \frac{d}{dt} \int_{S^1} \frac{1}{2} (\tilde{v} - v)^2 + \frac{\gamma}{2} (\tilde{u}_x - u_x)^2 dx \\ \leq C \int_{S^1} \frac{1}{2} (\tilde{v} - v)^2 + \frac{\gamma}{2} (\tilde{u}_x - u_x)^2 dx + \int_{S^1} E^2 dx. \end{aligned}$$

Error estimate

Applying Gronwall's Lemma and triangle inequality gives

Theorem (JG, Makridakis, Pryer '14)

Let (u_h, v_h) denote the solution of the semi-discrete dG scheme. Let (u, v) be a weak solution of (vdW) with

$$u \in C^1((0, T), L^2(S^1)) \cap C^0([0, T], H^3(S^1))$$

$$v \in C^1((0, T), L^2(S^1)) \cap C^0([0, T], H^1(S^1))$$

Then, there is $C > 0$ such that,

$$\|v_h(t, \cdot) - v(t, \cdot)\|_{L^2(S^1)}^2 + |u_h(t, \cdot) - u(t, \cdot)|_{dG}^2 \leq C(E_1 + E_2 + E_3),$$

where $E_1 := \left(\|\tilde{v}(0, \cdot) - v_0\|_{L^2(S^1)}^2 + |\tilde{u}(0, \cdot) - u_0|_{H^1(S^1)}^2 \right) \exp(Ct)$

$$E_2 := \int_0^t \int_{S^1} E^2 dx dt \cdot \exp(Ct)$$

$$E_3 := \|\tilde{v}(t, \cdot) - v_h(t, \cdot)\|_{L^2(S^1)}^2 + |\tilde{u}(t, \cdot) - u_h(t, \cdot)|_{dG}^2.$$

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Note that E_1, E_2, E_3 can be explicitly bounded in terms of u_h, v_h, τ_h .

Summary and outlook

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- New stability framework for regularized hyperbolic-elliptic problems.
- Derived an a posteriori error estimate.
- Estimate depends sensitively on γ , blows up for $\gamma \rightarrow 0$.
- Stability framework can also be used for model convergence.

Outlook

- Numerical experiments.
- Extension to fully discrete scheme.
- Extension to several space dimensions.
- Including viscosity.

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Thank you for your attention!